

Figure 14.14

Figure 14.14 shows a one-dimensional heat flow problem. The bottom end of a laterally insulated unit rod is immersed in a water solution at a fixed reference temperature. The temperature at the top end is also fixed at the same reference temperature. The *initial boundary value problem* (IBVP) of the heat flow problem is as follows:

| PDE | $\mathbf{u}_{t} = \alpha^{2} \mathbf{u}_{xx}$ | 0 < x < 1 | $0 < t < \infty$ |
|-----|---|-------------------|------------------|
| BCs | u (0,t) = 0 | 0 < t < ∞ | homogeneous BCs |
| | $u_x(1,t) + hu(1,t) = 0$ | | |
| IC | u (x,0) = x | $0 \leq x \leq 1$ | |

Determine the function **u** (**x**,**t**) by the **separation of variables method**.

where **u**(**x**,**t**) represents the temperature at some point **x** along the rod and at some point in time, **t**.

 $u_t = \delta u / \delta t;$ $u_{xx} = \delta^2 u / \delta x^2;$ h = heat exchange coefficient. α^2 = diffusivity (cm²/sec)PDE (partial differential equation)BCs (boundary conditions)IC (initial condition).

In Search Of The Function u (x,t) By The Separation Of Variables Method.

Assertion 1: there exists functions $X_n(x)$ and $T_n(t)$ such that:

$$u_n(x,t) = X_n(x) T_n(t)$$
 -----(1) (called fundamental solutions)

Assertion 2: the identity of u(x,t) is the same as the identity of the infinite sum of $u_n(x,t)$ that satisfies the given **IBVP**. That is:

$$\mathbf{u}(\mathbf{x},\mathbf{t}) = \sum_{n=1}^{\infty} \mathbf{A}_n \mathbf{u}_n(\mathbf{x},\mathbf{t}) = \sum_{n=1}^{\infty} \mathbf{A}_n \mathbf{X}_n(\mathbf{x}) \mathbf{T}_n(\mathbf{t})$$
-----(2) (if IBVP is satisfied)

Separating Variables:

$$\mathbf{u}(\mathbf{x},\mathbf{t}) = \mathbf{X}(\mathbf{x})\mathbf{T}(\mathbf{t})$$
 and $\mathbf{u}_{t} = \alpha^{2}\mathbf{u}_{xx}$

Implies:

So.

$$X(x)T'(t) = \alpha^2 X''(x)T(t)$$
-----(3)

where T'(t) = $\delta T/\delta t = u_t = \delta u/\delta t$; and X" = $\delta^2 X/\delta x^2 = u_{xx} = \delta^2 u/\delta x^2$

Dividing equation (3) by $\alpha^2 X(x)T(t)$, we have :

 $X(x)T'(t) / \alpha^{2}X(x)T(t) = \alpha^{2}X''(x)T(t) / \alpha^{2}X(x)T(t)$ $T'(t) / \alpha^{2}T(t) = X''(x) / X(x) - \dots - (4)$

The left hand side of equation (4) depends only on t and the right hand side depends only on x. Since x and t are independent, equation (4) implies that:

and
$$\begin{aligned} \mathbf{T}'(\mathbf{t})/\alpha^{2}\mathbf{T}(\mathbf{t}) &= \mu & \text{(where } \mu \text{ is the separation constant)} \\ \mathbf{X}''(\mathbf{x})/\mathbf{X}(\mathbf{x}) &= \mu \\ \text{So,} & \mathbf{T}'(\mathbf{t}) - \mu \alpha^{2}\mathbf{T}(\mathbf{t}) &= \mathbf{0} \text{-------(5)} \\ \text{and} & \mathbf{X}''(\mathbf{x}) - \mu \mathbf{X}(\mathbf{x}) &= \mathbf{0} \text{-------(6)} \end{aligned}$$

Equations (5) and (6) separate the variables and reduce the PDE to two ODEs.

T'(t) -
$$\mu \alpha^2 T(t) = 0$$
-----(5)
X''(x) - $\mu X(x) = 0$ -----(6)

 $\mu < 0$ is the domain of μ for which equations (5) and (6) are meaningful. If $\mu > 0$, u(x,t) = X(x)T(t) tends to infinity. If $\mu = 0$, u(x,t) = 0. μ is set equal to $-\lambda^2$ for $\mu < 0$. So, equations (5) and (6) become:

$$T'(t) + \lambda^2 \alpha^2 T(t) = 0$$

$$X''(x) + \lambda^2 X(x) = 0$$
(8)

The solutions for equations (7) and (8) are as follows:

So,
$$\mathbf{u}(\mathbf{x},t) = \mathbf{X}(\mathbf{x})\mathbf{T}(t) = e^{-(\lambda \alpha)^2 t} [A \sin(\lambda \mathbf{x}) + B \cos(\lambda \mathbf{x})]$$
------(11)

satisfies the PDE, $\mathbf{u}_t = \boldsymbol{\alpha}^2 \, \mathbf{u}_{xx}$ $\mathbf{0} < \mathbf{x} < \mathbf{1}$ $\mathbf{0} < \mathbf{t} < \infty$

for any λ and any A and B.

There are infinitely many **u** (**x**,**t**), as expressed in equation (11) that satisfy the PDE. We now look for those that satisfy both the PDE and the boundary conditions (BCs):

$$u(0,t) = 0$$

 $u_x(1,t) + hu(1,t) = 0$

So, substituting $e^{-(\lambda \alpha)^2 t} [A \sin(\lambda x) + B \cos(\lambda x)]$ into the BCs, we have:

$$Be^{-(\lambda \alpha)^{2}t} = 0 \implies B = 0$$
$$A\lambda e^{-(\lambda \alpha)^{2}t} \cos \lambda + hAe^{-(\lambda \alpha)^{2}t} \sin \lambda = 0$$
$$\tan \lambda = -\lambda/h - \dots - (12)$$

So,

$$\tan \lambda = -\lambda/h -----(12)$$

The values of λ for a given value of **h** (can be computed numerically with the aid of a computer) for which equation (12) is satisfied are called the eigenvalues of the boundary-value problem:

$$X''(x) + \lambda^2 X(x) = 0 -----(13)$$
$$X(0) = 0 -----(14)$$
$$X(1) + hX(1) = 0 -----(15)$$

These **eigenvalues** are the values of λ for which there exists a *nonzero solution* for the boundary-value problem. The solutions of the boundary-value problem derived from the eigenvalues λ_n are called the **eigenfunctions**, $X_n(x)$. For this boundary-value problem (equations 13 thru 15):

$$X_n(x) = \sin(\lambda_n x)$$

So, the infinite number of **fundamental functions** can be expressed as follows:

$$u_n(x,t) = X_n(x) T_n(t) = e^{-(\lambda_n \alpha)^2 t} \sin(\lambda_n x)$$
 -----(16)

Each of these functions satisfy the PDE and the BCs. Their sum such that the initial condition IC is satisfied is the identity of u(x,t).

So,
$$u(x, t) = \sum_{n=1}^{\infty} a_n e^{-(\lambda_n \alpha)^2 t} \sin(\lambda_n x) \quad ------(17)$$

such that the initial condition (IC), u(x, 0) = 0 is satisfied. That is:

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$$u(x, 0) = x = \sum_{n=1}^{\infty} a_n \sin(\lambda_n x)$$
 -----(18)

$$u(x, 0) = x = \sum_{n=1}^{\infty} a_n \sin(\lambda_n x)$$
 ------(18)

The constants \mathbf{a}_n in the eigenfunction expansion (equation 18) can be determined by multiplying each side of equation (18) by **sin** ($\lambda_m \mathbf{x}$) and integrating x from 0 to 1:

$$\int_{0}^{1} x \sin (\lambda_{m} x) dx = \sum_{n=1}^{\infty} a_{n} \int_{0}^{1} \sin (\lambda_{n} x) \sin (\lambda_{m} x) dx - \dots (19)$$

$$= 1$$
let x = ξ ; then, dx/d ξ = 1

So, equation (19) becomes:

$$\int_{0}^{1} \xi \sin (\lambda_{m} \xi) d\xi = \sum_{n=1}^{\infty} a_{n} \int_{0}^{1} \sin (\lambda_{n} \xi) \sin (\lambda_{m} \xi) d\xi -----(20)$$

$$= a_{m} \int_{0}^{1} \sin^{2} (\lambda_{m} \xi) d\xi$$

$$= a_{m} (\lambda_{m} - \sin \lambda_{m} \cos \lambda_{m})/2\lambda_{m}$$

$$a_{m} = 2\lambda_{m}/(\lambda_{m} - \sin \lambda_{m} \cos \lambda_{m}) \int_{0}^{1} \xi \sin (\lambda_{m} \xi) d\xi$$

Changing notation to *n*, we have:

So,

$$\mathbf{a}_{n} = 2\lambda_{n}/(\lambda_{n} - \sin \lambda_{n} \cos \lambda_{n}) \int_{0}^{1} \xi \sin (\lambda_{n} \xi) \, d\xi - \cdots - (21)$$

So, the solution to the IBVP problem is:

So,
$$\mathbf{u}(\mathbf{x}, \mathbf{t}) = \sum_{n=1}^{\infty} \mathbf{a}_n e^{-(\lambda_n \alpha)^2 t} \sin(\lambda_n \mathbf{x})$$
 ------(22)

where the constants \mathbf{a}_n are calculated from equation (21). Separation method is valid only for homogeneous IBVP. Other methods are used to solve non-homogeneous IBVP.

The String: S₇P₂A₂₁ (Identity – Physical Properties).

The Pj Problem of interest is of type i*dentity*. All problems of mathematical modeling are identity problems because the problems seek the mathematical structure of the physical problem being modeled.